

# Terahertz Imaging Systems With Aperture Synthesis Techniques

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**Abstract**—This paper presents the research and development of two terahertz imaging systems based on photonic and electronic principles, respectively. As part of this study, a survey of ongoing research in the field of terahertz imaging is provided focusing on security applications. Existing terahertz imaging systems are reviewed in terms of the employed architecture and data processing strategies. Active multichannel measurement method is found to be promising for real-time applications among the various terahertz imaging techniques and is chosen as a basis for the imaging instruments presented in this paper. An active system operation allows for a wide dynamic range, which is important for image quality. The described instruments employ a multichannel high-sensitivity heterodyne architecture and aperture filling techniques, with close to real-time image acquisition time. In the case of the photonic imaging system, mechanical scanning is completely obsolete. We show 2-D images of simulated 3-D image data for both systems. The reconstruction algorithms are suitable for 3-D real-time operation, only limited by mechanical scanning.

**Index Terms**—Femtosecond systems, photonic detectors, photonic systems, submillimeter-wave imaging, submillimeter-wave technology, synthetic aperture imaging, synthetic aperture radar, terahertz imaging.

## I. INTRODUCTION

THE MAJOR driving force for development of terahertz imaging systems today originates from security applications and, in particular, stand-off imaging of persons and hidden objects for security enforcement, including illicit drug and explosives detection. Bonding flaw and defect detection

at stand-off distances for space applications requires imaging systems with similar performance as those for security applications. Hence, system concepts can be discussed jointly for both applications. By terahertz, we mean a frequency range covering 300–3000 GHz, the lower part of which (100–1000 GHz) is sometimes referred to as the submillimeter-wave frequency region.

Conventional imaging systems utilize optical and X-ray imaging for stand-off detection for security and space applications [1]–[3], although a number of millimeter-wave, submillimeter-wave, and terahertz imaging systems have been presented [4]–[18]. In general, discussion on imaging systems can be broken down into passive imaging systems utilizing natural radiation or reflection and active imaging systems employing a dedicated source or an array of sources. Several recent overviews have been published on this subject [12], [16], [17], [19]–[21]. Another distinction on system level is time-domain versus frequency-domain operation.

This paper has two goals. The first one is to review imaging systems and algorithms at millimeter-wave and terahertz frequencies, subsequently called terahertz cameras. The second goal is to present two new imaging systems being developed by the authors.

The review focuses on active systems geared towards stand-off detection and fast image processing for nondestructive testing and security applications. Active imaging systems are preferred in unfavorable environmental conditions [19]. The resolution and the scene size require either an array or a scanning imaging architecture [22], [23]. Small receiver arrays have recently been presented in [24] and [25].

We will review systems based on electronic and photonic principles and provide a short description of imaging techniques and algorithms reported in the literature. As to electronic imaging systems, we mean those systems relying entirely on electronic components, whereas photonic imaging systems are those where a terahertz signal is generated or received through transformation of an optical signal into the terahertz range. Applications demanding spectroscopic capabilities for biological and chemical substance recognition are feasible at frequencies beyond 500 GHz, which are realized in this paper using a femtosecond multichannel imaging system.

One scenario considered in this paper is illustrated schematically in Fig. 1 and consists of an active array of transmitters and receivers illuminating an object at stand-off distance. The receivers are placed in the center part of the array, while the transmitters are situated at larger distances. The major imaging task here is to detect an object at a distance and classify the object

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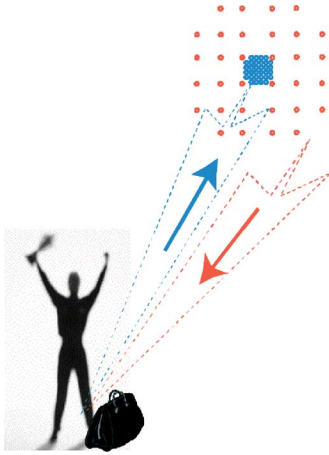


Fig. 1. Typical imaging scenario for stand-off detection using an active array of 32 transmitters (red, online version) and 32 receivers (blue, online version).

with regards to its potential threat or impact. One would like to detect all possible materials (i.e., solid, fluid, or gaseous) with many different properties (e.g., metal or dielectric). An imaging system in such a scenario should be able to detect the object with high spatial resolution and should classify this object with high success rate or low false alarm rate. The imaging should be close to real time in security applications with somewhat relaxed demands for industrial applications.

The terahertz camera considered here should operate at stand-off distances of a few meters to tens of meters, image man-sized scenes and produce images close to real time, and detect objects with a minimum size of a few centimeters. The bandwidth of the terahertz camera should therefore be large for good range resolution.

In the case of terahertz imaging systems based on photonic continuous-wave (CW) or pulsed techniques, the power is not sufficient to operate at such large distances, but the resolution is much better. This is due to the high operating frequencies. In this case, it is advantageous to operate a photonic terahertz imaging system at a closer distance in order to obtain submillimeter range and cross-range resolution. In fact, with the bandwidth of such systems, the range resolution is in the range of tens of micrometers.

## II. SURVEY OF ELECTRONIC IMAGING SYSTEMS

Active systems can operate in monostatic, bi-static, or multi-static mode. An advantage of active systems is the possibility for bi-static or multistatic measurements, providing information on the reflection and transmission properties of the object, which can be fully exploited in image processing algorithms. Passive systems operate only in an observation mode or with a broadband incoherent source, which then acts as a signal source [26]. However, powerful broadband noise sources are very difficult to realize. In principle, both passive and active imaging systems are feasible for the applications considered. Heterodyne millimeter-wave systems at room temperature require local oscillator (LO) signal sources, which can be used for object illumination.

Experiments demonstrated in literature [27]–[32] show that millimeter-wave frequencies ( $< 100$ – $300$  GHz) are sufficient

for good detection with high resolution while offering less complex systems and exhibiting better signal-to-noise ratios (SNRs) as compared with systems operating at higher frequencies. Another advantage of imaging systems at millimeter waves, as compared with terahertz systems, is the availability of integrated components and, in the near future, active monolithic integrated front-end modules, which would make multielement systems economically feasible. On the other hand, current experimental results reveal limited spectral features for substances such as explosives below 500 GHz and slightly lower frequencies for other materials [12].

### A. Passive Imaging Techniques

A passive imaging system is essentially a radiometer, which detects the radiation emitted from an object in view. One can employ either direct-detection focal plane arrays [33], [34] or a heterodyne receiver with a scanning system and single element receiver [35]–[40]. Passive multichip millimeter-wave imaging systems have been demonstrated in [41]–[46]. Direct-detection receivers have been using millimeter-wave monolithic microwave integrated circuit (MMIC) technology [10], [47] with good noise performance [4], [5], [8], [48]–[50], however, without dual-polarization and phase-detection capabilities. However, interferometric imaging requires the determination of the amplitude and phase of the detected signal [51], [52]. For adequate spatial resolution at large stand-off distances, apertures need to be large, i.e., of the order of 1-m diameter, and scanning should be employed [39], [53]–[55]. Other low-scan-rate mechanical systems have been developed in [56]–[58], and an alternative technology using a microbolometer cryogenic millimeter-wave focal plane array has been pursued in [33] and [58]. A stereoscopic passive imaging system has been presented in [59]. Recently, focal plane arrays at submillimeter-wave frequencies have been demonstrated operating at room temperature and using a standard CMOS process [14], [60], [61].

It is anticipated that, in the near future, MMIC technology will become the technology of choice from a cost and performance point of view, providing a natural path towards multichannel imaging systems.

### B. Active Imaging Techniques

Active imaging systems employ a transmitter to illuminate the scene and a receiver to collect the scattered energy, or multiples of either transmitters (Tx) or receivers (Rx), or both. Active imaging systems operating over a wide frequency range have been presented for security [3]–[18], [62]–[64] and biological [65], [66] applications.

In the case of arrays, one can imagine either focal plane or interferometric arrays including aperture filling techniques described below. Also, the receiver channels can use bolometric detection (direct detection) schemes or heterodyne receiver architectures. The latter is preferred in the case of active imaging systems due to the improved sensitivity, which could be two to three orders of magnitude better as compared with direct detection and provide micrometer-range resolution with large-bandwidth systems and utilizing the phase information.

Active imaging systems have the fundamental benefit of being independent of natural emission of the object and can therefore be successfully employed indoors and outdoors for concealed objects, even with thick cover materials, and are only limited by the available transmit power and receiver sensitivity. Another very important benefit is that the IF frequency bandwidth does not determine the overall system performance. Hence, IF frequency can be chosen as low as a few megahertz.

It should be emphasized that the emitted power is of the order of  $< 10$  mW and still well below any known health hazards [67] and therefore regarded as completely harmless. Thus, this technique could be ideal for large view scanning of populated areas.

A fully focused diffraction-limited 3-D image of a person or imaged target has been presented in [68] by scanning an inward-directed vertical array around the person or imaged target. The array achieved a full  $360^\circ$  mechanical scan in 2–10 s with optimized illumination, which minimizes signal loss due to specular reflection away from the array.

Time-domain electronic imaging systems utilize a pulse generated by fast switching of diodes [69], [70] or transistors [71] with bandwidths approaching those of photonic systems. In this case, all energy is concentrated in the pulse itself, which then should exhibit a sufficient energy level at all frequencies. Terahertz time-domain systems utilizing windowed Fourier transformation in 3-D have been presented in [72] and [73], and wide-angle tomographic imaging in synthetic aperture radar (SAR) has been reported in [74] employing mappings to modify the convolution-back projection method of computed tomography. Reconstructions from circular-arc projections that are made using the algorithm are of the same high quality as those made from standard straight-line projections.

### C. Scanners

A number of mechanically scanned optical systems have been developed, such as the spherical scanning antenna, the Lewis scanner, the Schwarzschild scanner, and the Rotman lens scanner [75]. Also, there is the scanning spherical trireflector antenna [76], the two-mirror scanning system [77], and the millimeter-wave monopulse twist reflector scanning antenna [78]. The oscillating motion of the mirrors used in these later designs is too slow to produce real-time scanning. A rotating polygon scanner [79], however, can scan at rates up to 50 kHz with scan angles approaching  $180^\circ$ . A high-speed scanning mirror [80] has formed the basis for a millimeter-wave line-scan system for use in airborne applications. Other imagers [81] achieve a raster scan by motion of both the antenna and receivers.

None of the systems described above are considered adequate for low-cost real-time imaging systems. Most commercial imaging systems employ some kind of scanners, with the exception of a recent electronically scanning system presented in [55] and [82].

It is therefore suggested not to use a scanner if a number of  $>16$  Tx and Rx channels can be afforded. Alternatively, one can use a linear array of elements for the imaging system and employ a simple scanner using a plane rotating mirror. This has

the advantage that optical defocusing and aberrations are kept to a minimum.

## III. SURVEY OF PHOTONIC IMAGING SYSTEMS

Photonic terahertz imaging systems have predominantly been employing time-of-flight measurements by raster scanning the imaging object. Recently, more advanced proposals for terahertz synthetic aperture and interferometric imaging methods have been proposed by several groups and published both in the scientific literature and filed as patents. A comprehensive review of imaging with terahertz radiation has recently been published in [11], [12], and [85]–[88].

In practice, until today, most systems employ one Tx/Rx pair by moving the transmitter and/or receiver around the imaging object, acquiring data from various positions and subsequently reconstructing an image. Having more than one transmitter and/or receiver at different positions enables illumination of an object from more than one location and, hence, employs tomography. The advent of the fiber-coupled photoconductive emitters/receivers made this a viable method [89]–[92]. Below, we will demonstrate a possible implementation of a multi-channel imaging system using fiber-coupled sensor heads.

Several research groups have explored the analogy of X-ray computer tomography (CT) with terahertz radiation and pioneered by Zhang *et al.* [91]. In contrast to X-ray CT, terahertz CT resolves both the amplitude and the phase information of the scattering object. As a consequence, the terahertz CT image contains more information than the X-ray CT about the target, such as the frequency-resolved refractive index. The wide-aperture reflection tomography allows tomographic reconstruction of a series of slices measured at different view angles [94]. This technique works best with strong reflectors such as metals. The algorithms used for the reconstruction of the images are in general filtered back-projection algorithms.

## IV. SURVEY OF IMAGING ALGORITHMS FOR TERAHERTZ AND MILLIMETER-WAVE SYSTEMS

Focusing the raw terahertz data of a terahertz imaging system for security and space applications is a challenging task due to the required real-time processing, 3-D mapping, the comparable object and antenna aperture size with the stand-off distance, the relatively large bandwidth reaching 1000 GHz, and the adoption of the bi-static configuration [95].

One review [96] covers the fundamental principles of radar imaging addressing aspects of through-wall radar imaging, radar detection of buried targets, tomography and detection of concealed weapons, and passive bi-static radar.

### A. Resolution

The theoretical range resolution  $\rho$  depends on the bandwidth of the transmitted signal  $B$  [98]

$$\rho = \frac{c}{2B}. \quad (1)$$

A 2-cm resolution, for instance, requires a 7.5-GHz bandwidth without additional phase processing.

In the cross-range direction, a passive diffraction-limited system would offer an angular resolution of

$$\alpha = \frac{\lambda}{D} \quad (2)$$

where  $\lambda$  is the wavelength of the incident radiation and  $D$  is the diameter of the aperture. The spatial cross-range resolution is found by multiplying  $\alpha$  (in radians) with the range to the object. For an active monostatic system, the angular resolution is half of that in (2) [98], whereas, for a multistatic system like that in Fig. 1, (2) holds true for an object in the far field, as the effective size of the aperture is half the size of the transmitter aperture  $D$ .

A multistatic radar transmits sequentially on each of the  $N_{TX}$  TX antennas and receives the scattered signal simultaneously on each of the  $N_{RX}$  RX antennas. For a TX, RX element pair, the round-trip equirange surfaces are ellipsoids, which in the far field approximate spheres centered at the midpoint between the TX and RX antennas.

The angular field of view (FOV) is determined by the smaller of two angles: the antenna element beam width and the angular ambiguity spacing (grating lobe spacing). If the beam width exceeds the ambiguity spacing, the imaged object must be sufficiently small to fit in between two ambiguities.

The angular ambiguity spacing  $\theta$  depends on the effective element spacing  $d$

$$\theta = \frac{\lambda}{d} = \frac{\lambda N}{D} \quad (3)$$

where  $N$  is the effective number of elements. Hence, large objects call for closely spaced elements. Equation (3) applies to a multistatic system, for which the ambiguity spacing is twice that of a monostatic radar system [99], as its effective elements spacing is half of the physical element spacing.

The maximum number of cross-range pixels that can be resolved  $N_{cr}$  is the ratio of the ambiguity spacing to the angular resolution

$$N_{cr} = \frac{\theta}{\alpha} \quad (4)$$

which, according to (2) and (3), equals the effective number of array elements  $N$ . As an example, an imager with a 1-m aperture operating at a wavelength of 1 mm and covering a FOV of  $12^\circ$  would have an angular resolution of  $0.06^\circ$ , and it would call for at least 200 elements and produce the same number of uncorrelated pixels. For comparison, a typical thermal imager operating at  $10 \mu\text{m}$  and having an objective diameter of 50 mm and an FOV of  $12^\circ$  has a theoretical number of pixels of about 1000. In practice, aberrations introduced by the detector array lower this value to be closer to 500.

### B. Aperture Filling

In order to avoid grating lobes while obtaining the best possible resolution and the maximum unambiguous scene size with a given number of antenna elements, the  $N_{TX}N_{RX}$  midpoints must be equidistantly spaced. This can be obtained with a TX antenna element spacing of

$$\Delta_{TX} = N_{RX}\Delta_{RX} \quad (5)$$

where  $\Delta_{RX}$  is the RX antenna element spacing. (Here and in the following, the TX elements are the more widely spaced elements, but TX and RX can obviously be interchanged.) It is noted that (5) implies no redundancy in the sense that no TX–RX midpoints coincide. The overall size of the antenna array is  $N_{TX}\Delta_{TX} = N_{TX}N_{RX}\Delta_{RX}$ , which offers an angular resolution of

$$\alpha = \frac{\lambda}{N_{TX}N_{RX}\Delta_{RX}}. \quad (6)$$

The angular ambiguity separation is

$$\theta = \frac{\lambda}{\Delta_{RX}} = N_{TX}N_{RX}\alpha. \quad (7)$$

Hence, a multistatic radar with  $N_{TX} + N_{RX}$  antenna elements can resolve  $N_{TX}N_{RX}$  uncorrelated pixels, which is typically much more than the  $N_{TX} + N_{RX}$  pixels that can be resolved with a monostatic radar with the same number of elements. The same principle can be applied to a 2-D array [120].

It should be noted that the aperture filling outlined above is fundamentally different from that of synthetic aperture radiometers as used in radio astronomy [122] and in Earth observation [123]. Such correlation radiometers accomplish the aperture filling by cross correlating the signals received by any two pair of receiving elements. In order to fill the aperture, all multiples of the minimum baseline (i.e., the minimum element spacing) must be represented. Low-redundancy configurations are known for 1-D arrays, e.g., [99], but not for 2-D arrays [122]. For instance, the Y-shaped array adopted by the SMOS system does imply some redundancy.

Aperture filling by means of cross correlation cannot be applied to radar imaging of distributed targets because of the nonzero cross correlation of the signals scattered by different scatterers. However, for imaging of objects consisting of a few reflectors on a nonreflecting background, the principle may be applicable.

### C. Range Focusing

A large time-bandwidth product of the transmitted signal is necessary in order to obtain a good SNR (long pulse) simultaneously with good resolution (large bandwidth). Using a delayed replica of the transmitted signal for demodulation, a target at a given range will result in a CW signal with a frequency proportional to the range from some reference range in deramp systems employing linear frequency-modulated (FM) signals. Hence, the focused range response can be obtained using a fast Fourier transform (FFT). The offset of the reference signal can be chosen so that a target at the range window center will exhibit a frequency of 0 Hz (with I/Q demodulation) or some IF frequency sufficiently low to allow digitization. For a scene with a range extent of  $\Delta r$ , the required IF bandwidth  $B_{IF}$  (and analog-to-digital converter (ADC) bandwidth) after analog deramp is a fraction of the modulation bandwidth  $B_p$

$$B_{IF} = \frac{2\Delta r}{c\tau_p} B_p = \frac{\tau_f}{\tau_p} B_p \quad (\text{Deramp, single transmitter}). \quad (8)$$

This simplifies the baseband system when the temporal range window  $\tau_f$  is shorter than the uncompressed pulse length  $\tau_p$ .

The minimum pulse length that must be recorded to focus the range window is  $\tau_r = \tau_p + \tau_f$ . For a system with multiple transmitters, it is desirable to transmit simultaneously on all transmitters in the array, separating the transmitted signals in frequency. The frequency separation  $\Delta f_{\text{TX}}$  should be small compared with the pulse bandwidth  $B_p$  (to allow SAR focusing as well as digitization) but larger than the  $B_{\text{IF}}$  given in (8) (to separate transmitters in the recorded data). The required IF and ADC bandwidth are then

$$B_{\text{IF}} \geq N_{\text{TX}} \Delta f_{\text{TX}} \geq N_{\text{TX}} \frac{\tau_f}{\tau_p} B_p \quad (\text{multiple transmitters}). \quad (9)$$

The transmitters are easily separated after digitization, as the frequency offset is converted to a range offset by the range-focusing FFT. The ambiguous range is then determined by the transmitter frequency separation

$$r_{\text{amb}} = \frac{c \Delta f_{\text{TX}} \tau_p}{2 B_p}. \quad (10)$$

#### D. Cross-Range Focusing

The matched filter algorithm correlates raw data with the system point target response [97]. The shift-variance of the correlation step means that FFT-based filtering cannot be applied, resulting in a very large number of arithmetic operations per pixel. However, matched filtering is an optimum algorithm in that it does not exploit any approximations. Adjacent resolution-cell leakage (sidelobes) can be suppressed by amplitude weighting of the filter coefficients. Matched filtering is a very flexible algorithm, allowing for 2-D or 3-D focusing, without imposing any constraints on the placement of the antenna elements.

The mapping geometry of typical terahertz imaging systems is similar to that of spotlight SAR in the sense that a geometrically confined scene/object is illuminated by transmitter antenna elements having an angular separation with respect to the object. In its simplest form, a spotlight SAR system is a monostatic radar with the object in the far field of the synthesized antenna. In this case, the simple rectangular algorithm [98] can be employed. Unfortunately, this is hardly applicable to terahertz imaging data, because the object is in the near field of the antenna array and a scatterer is not observed in the same range cell by all array elements. The Polar Format Algorithm [98]–[101] compensates to some extent for this phenomenon and can therefore be applied to scenes much larger than those covered by the rectangular format algorithm. Different formats have been studied in [102] including the close-range scenarios.

The SAR imaging process implies that the complex signals from all scatterers at the same round-trip range add up and form a sample of the (range-compressed) receive pulse. The result of the projection can be expressed as a sample of the Radon transform, which is the integral of the 3-D scene reflectivity over a plane [103]. The scene reflectivity can then be recovered by an inversion of the Radon transform for circular and spherical geometries [104] and for multistatic scenarios [105], [106]. A different approach has been proposed in [107], which utilizes the reversible transform between time delay and target boundary and has been applied to moving persons in [108],

where moving motion of the persons was employed instead of scanning apertures.

One way of implementing the inverse Radon transform is the back projection algorithm (BPA), which is a spatial domain technique and, hence, somewhat computationally intensive. For every pulse and for every pixel in the output image, an interpolated sample must be computed and accumulated. “Fast” BPAs do exist [109]–[112], including near-field scenarios, with improvements in speed by a factor of 100 compared with direct back-projection in 2-D [111]. The advantage of these algorithms is that they provide some of the flexibility and robustness of a time-domain algorithm but exhibit a computational performance in parity with fast transform algorithms, e.g., the Fourier–Hankel or range migration algorithms (RMAs). Fast BPA for simple bi-static geometries have been proposed in [113] with a tradeoff in image quality. The BPA has also many attractive properties in common with the matched filtering: great flexibility, little constraint on the placement of the antenna elements, support of a large relative bandwidth, and a large relative aperture or object. It can be shown that with a perfect interpolator, the BPA is an exact inversion of the linear aperture SAR problem [111].

The RMA [114]–[116] is known under a number of different names, e.g.,  $\omega$ -k algorithm, wave equation technique, and seismic migration technique. For the center scatterer, the RMA is equivalent to the matched filter algorithm, and it is optimum in the same sense. However, the RMA implements the matched filtering in the spectral domain and, consequently, applies the same filter to the entire scene. For scatterers displaced from the scene center, a correction with the so-called Stolt transformation should be employed according to  $K_Y = \sqrt{K_R^2 - K_X^2}$ , which is a transformation from the range wavenumber in signal space  $K_R$  to the range wavenumber in object space  $K_Y$ .  $K_X$  is the cross-range wavenumber. The Stolt transformation corrects in a perfect way for the residual range migration and phase errors at all frequencies, but it is very sensitive to processing artifacts and offers little flexibility in the placement of the antenna elements. The RMA is well suited for wideband systems and for near-field systems with large apertures or large objects compared with the stand-off distance. A number of algorithms similar to RMA have been proposed, such as the chirp scaling algorithm [117], [118], where the final range inverse FFT is replaced by a single nonuniform FFT. Nonuniform FFT based projector/back projector with  $O(N^2 \log(N))$  complexity has been presented in [119]. A comparison of the BPA and the RMA is provided in [121].

For real-time or close to real-time applications, it seems that the BPA is a good compromise between performance, flexibility, and speed.

#### V. PHOTONIC TERAHERTZ MULTIELEMENT IMAGING SYSTEM REALIZATION

As discussed above, most photonic imaging systems today utilize a single element with the mechanical scanning, resulting in long acquisition times. The photonic terahertz imaging system presented here is based on an advanced dual-femtosecond fiber laser with an electronic delay stage unit, an all-fiber ultrashort pulse distribution system with dispersion

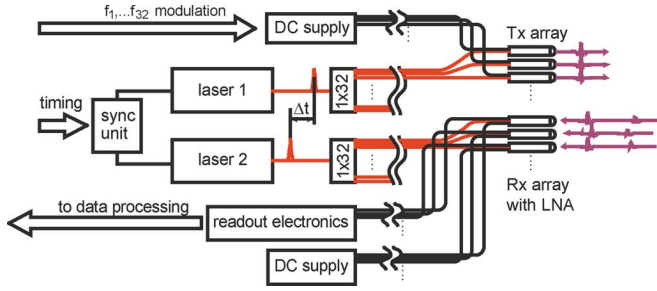


Fig. 2. Complete photonic terahertz imaging system, including electronic components. The thick arrows indicate communication with the data acquisition back-end.

control [135], and a 2-D array of 32 individual terahertz emitters and 32 individual terahertz detectors. The advantage of an all-fiber design is a system that is very stable with respect to external vibrations and temperature fluctuations, without need for further alignment. In addition, we have developed ultra low-noise amplifiers, directly integrated into the individual terahertz receivers. Fig. 2 shows an outline of the principal components of the photonic imaging system. The first experimental results for this system are described below.

#### A. Laser System and Pulse Delivery

We have chosen a dual-femtosecond fiber-based laser system from Toptica Photonics AG with approximately 4.5 nJ per laser, 0.9-ps laser pulses around 1550 nm with a 90-MHz pulse repetition rate. The two femtosecond lasers are locked to the same repetition rate by active control, and the relative time delay can be adjusted electronically between 0.0–1.0 ns. This enables electronically controlled optical sampling (ECOPS). All laser pulses were measured by frequency-resolved optical gating (FROG) [134], and the retrieved pulses were compared with the measured intensity autocorrelation. The time jitter of the two synchronized lasers was characterized by measuring the intensity cross correlation.

It is a nontrivial task to distribute the high-power laser pulses to the 32-THz emitters and detectors. We have designed a fiber link, indicated in Fig. 3, in which the laser pulses are first stretched, then split into 32 channels, and finally compressed such that they are around 100 fs when they reach the emitters and detectors in order to obtain efficient terahertz generation. In this setup, the dominant nonlinear effect is self-phase modulation, which adds an intensity-dependent phase shift to the pulse. We first stretch the pulse for as large a fraction of the fiber link as possible using a dispersion-compensating fiber (DCF) module from OFS Fitel, resulting in a pulsewidth of around 2.2 ps. The stretched laser pulse is then split into 32 equal portions in a  $1 \times 32$  power splitter from Ignis Photonix capable of handling average optical powers  $< 440$  mW. The insertion loss was around 1.5 dB. The normal dispersion of the DCF is compensated by the anomalous dispersion of the single-mode fiber (SMF) after the splitter. The optimal SMF fiber length of 4.6 m was found in a cut-back experiment, resulting in minimum pulse duration of 88 fs, as illustrated in Fig. 4.

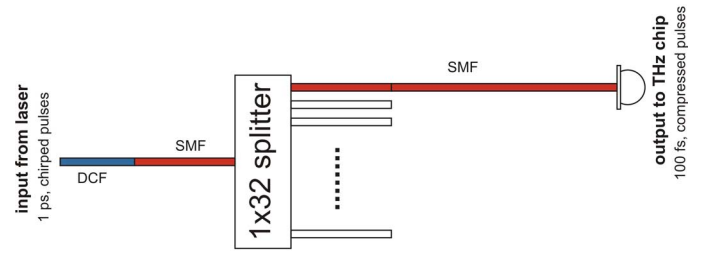


Fig. 3. Schematic of the fiber link. The input laser pulse is stretched in a 50-cm dispersion-compensating fiber (DCF) module with single-mode fiber (SMF) pigtailed. The DCF module is spliced onto a  $1 \times 32$  splitter, and each of these channels is spliced onto a piece of SMF (around 3 m) for pulse compression.

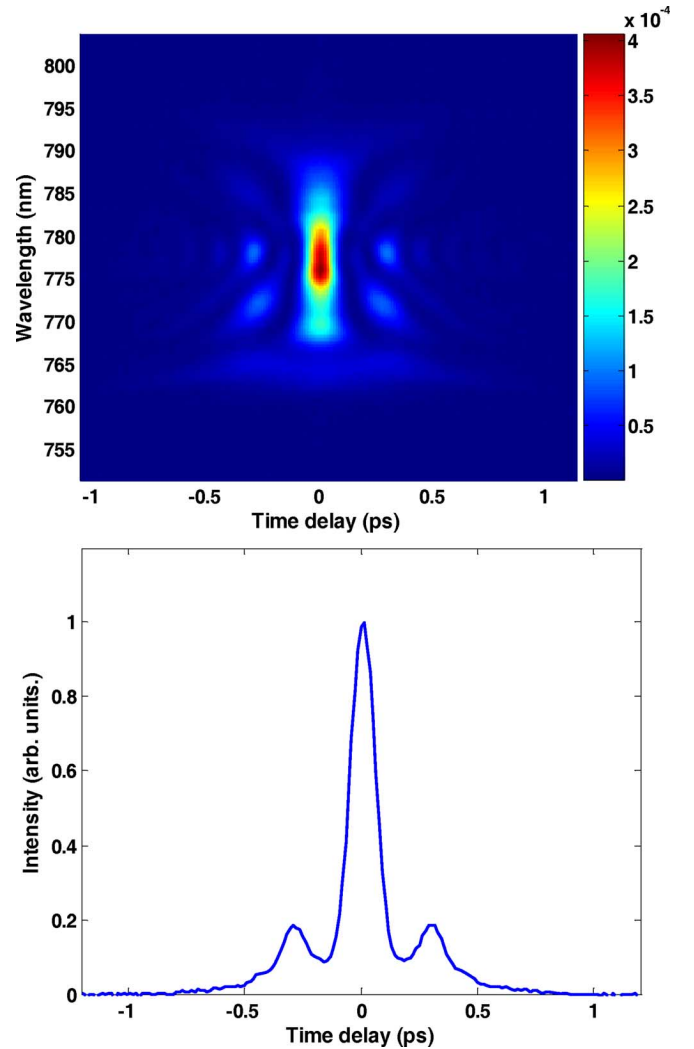


Fig. 4. Measurements of the laser pulse at the end of the fiber link. Top: Measured FROG trace. The FWHM of the retrieved pulse is 88 fs. Bottom: Intensity autocorrelation of the pulse. The FWHM of the autocorrelation is 131 fs.

#### B. Terahertz Emitters and Detectors

Both terahertz emitters and receivers use photoconductive switches to respectively generate and receive terahertz radiation [133] using low-temperature-grown InGaAs layers [91] at 1550 nm, which are commercially available from Menlo Systems GmbH.



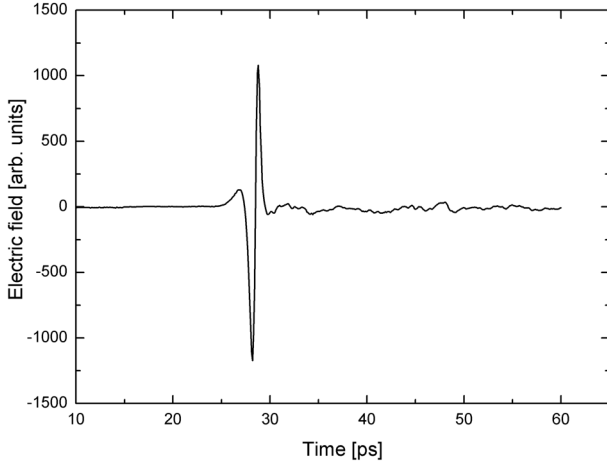


Fig. 5. Terahertz pulse generated and received with photoconductive switches from Menlo Systems fed by femtosecond laser pulses from the fiber link.

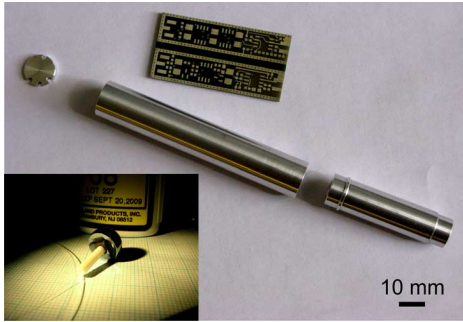


Fig. 6. Photograph of the parts of the terahertz sensor head. The diameter of the heads is 12 mm, and the total length is 130 mm. The silicon lenses have a diameter of 10 mm and are mounted into the rightmost part. The PCB shows the unpopulated current amplifier and fits into the tube shown in the center of the picture. Inset: The photoconductive antenna chip glued directly onto the silicon lens. The fiber link is butt-coupled into the photoconducting gap.

The terahertz radiation is collimated out of the photoconductor using hyper-hemispherical high-resistivity silicon lenses. Fig. 5 shows a femtosecond terahertz pulse generated and received with the photoconductive switches described above and fed by direct butt-coupling of the fiber.

The ultracompact design of the transmitter and receiver head integrates silicon-lens, photoconductive switch, and, respectively, bias supply (in the emitter) and low-noise current amplifier (in the receiver) in one single compact unit, as depicted in Fig. 6. The custom-designed LNA preamplifier exhibits a transimpedance gain of  $10^9$  V/A, an equivalent noise voltage and current of  $e_n \approx 6$  nV/sqrt(Hz),  $i_n \approx 1$  fA/sqrt(Hz), with a bandwidth of  $\sim 20$  kHz, and an input capacitance of  $\sim 3$  pF.

### C. Design of the 2-D Synthetic Aperture Array

The photonic system consists of a 2-D transmit/receive synthetic aperture array. The practical realization is depicted in Fig. 7. The 32 receivers are stacked in a  $6 \times 6$  grid (corners omitted for ease of manufacturing) in the inner part of the array and spaced by 13 mm. The transmitter units are placed in the outer part and spaced by 80 mm. All units are fitted into a 20

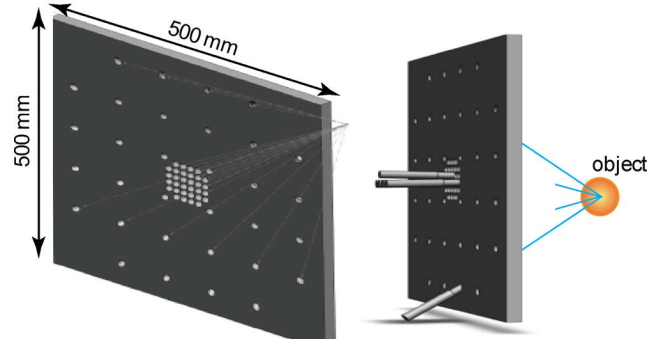


Fig. 7. 3-D rendering of the mounting structure for the photonic terahertz antenna array. Each element is directed towards the scene center by the structure. The receivers are placed in the inner array spaced by 13 mm and the transmitters in the outer array spaced by 80 mm.

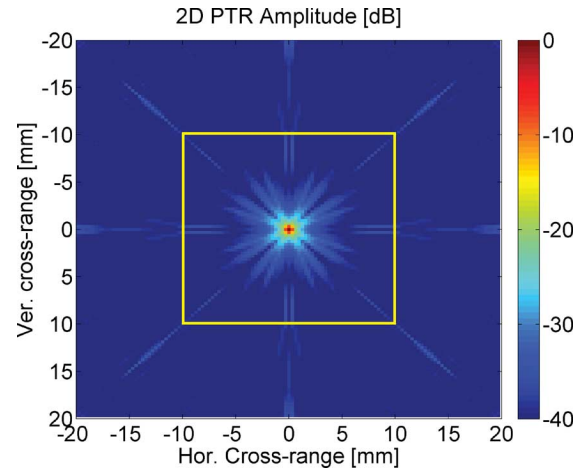


Fig. 8. Point target reconstruction simulation for the  $32 \times 32$  element photonic array with a bandwidth of 100–1000 GHz. Red square (online version) indicates desired cross-range scene extent (2 cm  $\times$  2 cm).

mm  $\times$  500 mm  $\times$  500 mm aluminum mounting structure and pointing towards a center of scene at 300 mm from the structure.

The photonic terahertz imaging system has not yet been finalized, and therefore simulation results are shown in this paper to demonstrate the performance of the imaging system. All simulations were performed for a Gaussian pulse system with a bandwidth of 900 GHz centered around a frequency of 550 GHz. A simulated point target reconstruction is shown in Fig. 8, with the slice shown taken at the range of the point target. The simulated cross-range 3-dB resolution is 0.4 mm [consistent with (2)], and the first ambiguities are centered approximately 14 mm from the point target [consistent with (3)]. The ambiguities are smeared due to the large bandwidth employed and are 35 dB below the peak. If a full volume (2 cm (range)  $\times$  4 cm  $\times$  4 cm) is reconstructed and integrated in range, the peak relative ambiguity level rises to around  $-20$  dB.

## VI. MULTIELEMENT ELECTRONIC IMAGING SYSTEM REALIZATION

The electronic imaging system combines a horizontal, linear array of elements with vertically scanning optics and will be capable of scanning a 1 m  $\times$  2 m scene at a distance  $> 7$ -m scene

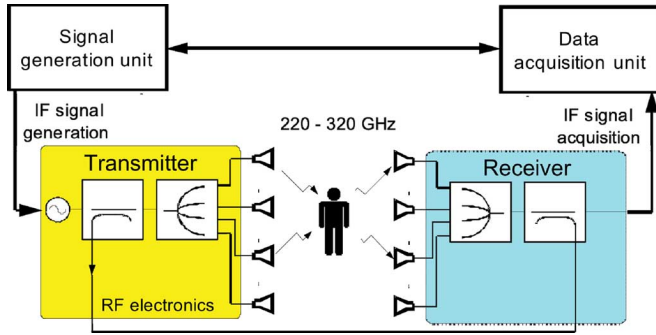


Fig. 9. Block diagram of the electronic imaging system. The component choice is given in the text.

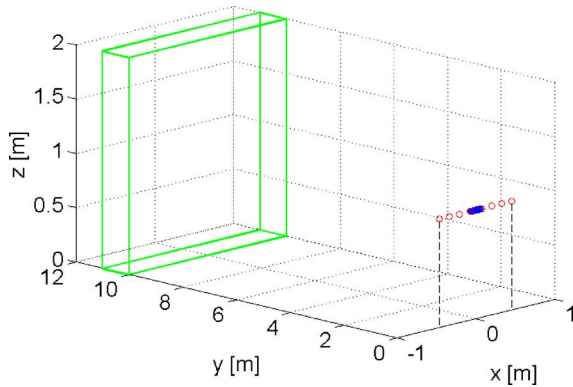


Fig. 10. Sketch of the scene (green, online version) and array configuration of the electronic imaging system with 16 receiving (blue, online version) and eight transmitting elements (red, online version) at a working distance of 10 m. The vertical scanning optics have been omitted for clarity.

with the according reflector dimensions. The image acquisition and processing is foreseen to be in real time, enabling imaging of moving objects. Using an image refresh rate of  $t_{\text{ref}} < 500$  ms and 100 lines in the vertical direction (2-cm resolution of a 2-m scene) leaves 5-ms/line acquisition time; this implies that the optical scanner has to move faster than this time. This puts high demands on the mechanical design and the data acquisition unit.

The suggested electronic imaging system features an array of Tx and Rx channels in multistatic configuration operating at 220–320 GHz (100-GHz bandwidth), a data acquisition and signal generation unit, as shown in Fig. 9. The delay line connecting the transmitter and the receiver is a fixed delay line employing a 50- $\Omega$  coaxial cable.

In general, front-end electronic components operate over a bandwidth typical for waveguide components. Wider bandwidth is currently very difficult to achieve in electronic systems [69]–[71], [84].

Wideband power dividers at the  $Ku$ -band have been developed to distribute the LO signals, which are then multiplied to the output frequency. The element spacing for the Tx and Rx units have been optimized for low redundancy in the array and large unambiguous scene, as discussed above. Currently, the system employs 16 Rx units and 8 Tx units with 8-mm receiver and 128-mm transmitter spacing at a stand off distance of 10 m, illustrated in Fig. 10.

One bottleneck in development of multichannel terahertz imaging systems is the data transfer of the acquired data.

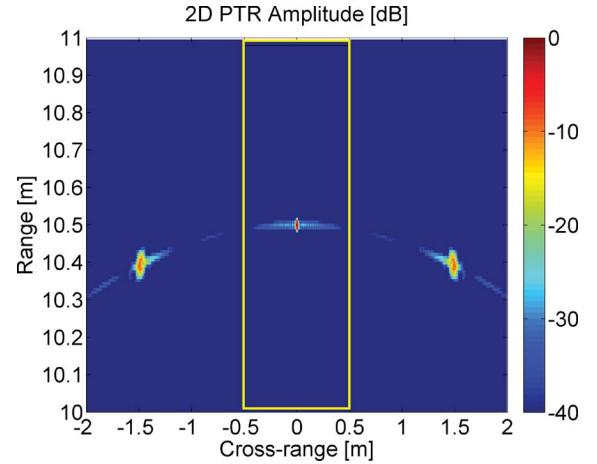


Fig. 11. Electronic imaging system simulated image reconstruction for a point target using the BPA. Red (online version) rectangle indicates desired scene size (1 m)—an extended scene was simulated to illustrate ambiguities.

For example, sampling the baseband image data from the imaging system with 30 MS/s in a 32-channel system with 2 bytes/sample and 10 images/s results in 1920-MB/s data stream, assuming an image acquisition time of 0.1 s/image. Such a huge data stream is difficult to support with standard commercial components utilizing, e.g., the PCIe bus. Also, the real-time computation of such large amounts of data is difficult to achieve using commercial off-the-shelf components. Using superheterodyne architectures alleviates this problem at the expense of larger complexity.

Synthetic aperture reconstruction for the configuration in Fig. 10 has been simulated assuming a point target at the center of the scene. The resulting  $x$ – $y$  slice ( $z = 1$  m) is shown in Fig. 11. Hamming weighting of the aperture elements was employed to suppress cross-range sidelobes, giving a measured resolution of 1.4 cm, which is slightly larger than expected from (3), which does not assume weighting. Ambiguities are 1.5 m away from the point target, consistent with (4), and well outside the desired 1-m scene width. Real-time reconstruction was demonstrated using a parallel implementation of the BPA on an NVIDIA GTX260 graphics processing unit (GPU). A  $1 \times 1$  m  $xy$ -slice of  $128 \times 128$  pixels could be focused in less than 2 ms.

## VII. CONCLUSION

This paper reviewed performance aspects of published terahertz imaging systems with special emphasis on security and object detection at stand-off distances. Literature results indicate that mechanical scanning represent a bottleneck in real-time imaging systems, whereas, for array imaging systems, one bottleneck is the data throughput. There is a clear tendency towards monolithic integration of front-ends for passive and active imaging systems. Imaging system parameters such as bandwidth, number of elements, element spacing, and antenna radiation pattern have been established for active electronic and photonic imaging systems, respectively. The discussion on most appropriate algorithms resulted in the choice of the BPA as the most suitable for real-time imaging system operation and good image quality.



Based on this discussion, an electronic and photonic terahertz imaging system has been demonstrated, and preliminary measured results for components of these systems have been shown. The overall system performance has been demonstrated using image reconstruction of point target data.

The photonic system comprises a fiber distribution system to the individual 32 Tx and 32 Rx units. Measured results indicate that 80-fs high-power pulses are obtained at the input to the respective terahertz antennas, which demonstrates state-of-the-art performance. Each terahertz antenna is integrated with state-of-the-art LNA. The Tx and Rx 2-D arrangement has been optimized for best image quality with a spatial resolution of 0.4 mm and an ambiguity level  $< -20$  dB. The image processing speed for this imaging system is  $< 1$  ms.

The electronic imaging system is based on a linear array operating in a frequency range of 220–320 GHz and employing an optical scanner for stand-off detection at  $> 7$  m. It is demonstrated that the image processing for this system is  $< 2$  ms with a spatial resolution of  $< 1.5$  cm. The development of this system is still under way, but simulations of the overall system performance demonstrate that ambiguities are 1.5 m away from the point target, as necessary for person scanning. The reconstruction of the full scene of  $1 \text{ m} \times 2 \text{ m}$  is limited by the scanning mirror movement.

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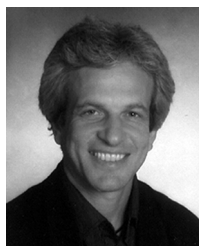
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